

Adaptive cytoskeletal responses to extracellular environment viscosity modulate cell migration

Zhongya Lin¹, Xindong Chen^{1,*}, and Xi-Qiao Feng^{1,2,*}

¹Institute of Biomechanics and Medical Engineering, AML, Department of Engineering Mechanics, Tsinghua University, Beijing 100084, China

²Mechano-X Institute, Tsinghua University, Beijing 100084, China

*E-mail: chenxin7@tsinghua.edu.cn, fengxq@tsinghua.edu.cn

ABSTRACT

Cell migration is a pivotal process in metastasis, allowing cancer cells to invade surrounding tissues and disseminate to distant organs. While extracellular environment (ECE) viscosity serves as a critical modulator of cell motility, its regulatory mechanisms remain unclear. This study presents a mechanobiological model to investigate how ECE viscosity modulates cancer cell migration by regulating some key processes, including actin polymerization, retrograde flow, and adhesion adaptations. Our results reveal a biphasic response: a moderate increase in ECE viscosity enhances actin filament network density and adhesion strength, thereby accelerating migration, whereas excessively high viscosity hinders movement due to too large mechanical resistance. Furthermore, we identify a short-term migration memory phenomenon, where cancer cells exposed to high viscosity environments retain elevated migration speeds after transitioning to low viscosity conditions. This memory effect is sustained by the continued assembly of cytoskeletal proteins such as actin monomers and Arp2/3. These analyses reveal an adaptive mechano-chemo-biological mechanism by which cancer cells integrate and respond to mechanical cues from their viscous environment to optimize migration, and advance the understanding of cancer cell migration in various tissue environments.

keywords: viscosity, actin filament dynamics, cell migration, short-term memory

¹ Introduction

Cell migration underpins numerous physiological and pathological phenomena. Examples include the migration of stem cells to establish the body plan and organ systems during the earliest stages of embryogenesis^{1,2}, the intricate movements of immune cells in response to infection^{3,4}, and the complex dynamics of cancer metastasis^{5,6}. Actin cytoskeleton is a primary force-generating structure that drives cell motility, mainly through dynamic assembly and disassembly processes at the leading edge of migrating cells^{7–9}. Actin filaments (AFs) rapidly polymerize to form structures such as lamellipodia and filopodia, which are essential for protrusion and adhesion to the extracellular environment^{10,11}. Cells explore their environment and adapt their subsequent responses through dynamic cytoskeletal processes. Therefore, the physical properties of the extracellular environment, including stiffness, topography, crosslinking, viscoelasticity, and porosity, significantly influence cellular behavior and actin dynamics^{12–17}. Among these factors, the role of viscosity has emerged as a critical but relatively understudied modulator of cell motility, actively shaping the biomechanical behavior of cells^{18–20}.

Elevated viscosity has been observed in various types of tumors, often due to the accumulation of cancer cells

14 and their secreted factors^{16,21}. Interestingly, recent studies have shown that increased ECE viscosity can enhance the
15 motility of certain cancer cell types^{20,22,23}, challenging the conventional expectation that higher resistance would
16 suppress cell movement. This counterintuitive finding suggests that cancer cells have evolved adaptive mechanisms
17 to navigate more efficiently through viscous environments, which may be driven by the alterations in actin dynamics,
18 ion channel activity, and cell-matrix interactions. Furthermore, the changes in ECE viscosity may lead to migration
19 memory, where cancer cells retain their migratory behavior after exposure to different viscosities²³. Mechanical
20 memory may significantly enhance their ability to survive and thrive in distant organs, whose mechanisms remain to
21 be elucidated²⁴⁻²⁶. Therefore, a thorough understanding of how ECE viscosity modulates cell migration is crucial
22 for decoding the mechanisms that drive tumor progression and metastasis.

23 Although the importance of actin dynamics in cell motility has been well recognized, the effects of varying ECE
24 viscosity on the organization and behavior of the AF network have not been thoroughly investigated. Furthermore,
25 while several recent works have addressed how viscosity influences actin dynamics, they overlooked the critical
26 feedback mechanisms between cellular components that lead to mechanical memory. To address this gap, we
27 establish a multiscale mechanobiological model that simulates the evolution of the AF network at the leading edge in
28 response to changes in ECE viscosity. We aim to elucidate how variations in viscosity influence actin polymerization,
29 retrograde flow, and cell-matrix adhesions, thereby characterizing the mechanistic influences of ECE viscosity
30 on cancer cell migration. Furthermore, we investigate how exposure to different viscosities can induce lasting
31 biomechanical adaptations, i.e., cytoskeleton-based short-term migration memory, in cancer cells, providing new
32 insights into their ability to optimize movement in response to different mechanical cues.

33 **Results**

34 **Mechanobiological model of cell migration**

35 In this paper, we focus on lamellipodia induced migration. Lamellipodia are characterized by a branched network of
36 AFs, primarily formed through the action of Arp2/3, which serves as a nucleation site for new filament assembly^{27,28}.
37 This protein promotes the formation of a dense network that drives the protrusion of the cell membrane. For a cancer
38 cell migrating in the viscous ECE, the driving force is provided by the branched AF network near the leading edge,
39 which generates sufficient propulsion for cell movement (Fig. 1a). The leading edge is subject to the propulsive
40 force of AF polymerization, the resistance of cell membrane tension and the ECE (Fig. 1b). At the steady state, the
41 force balance equation at the leading edge of a migrating cell is

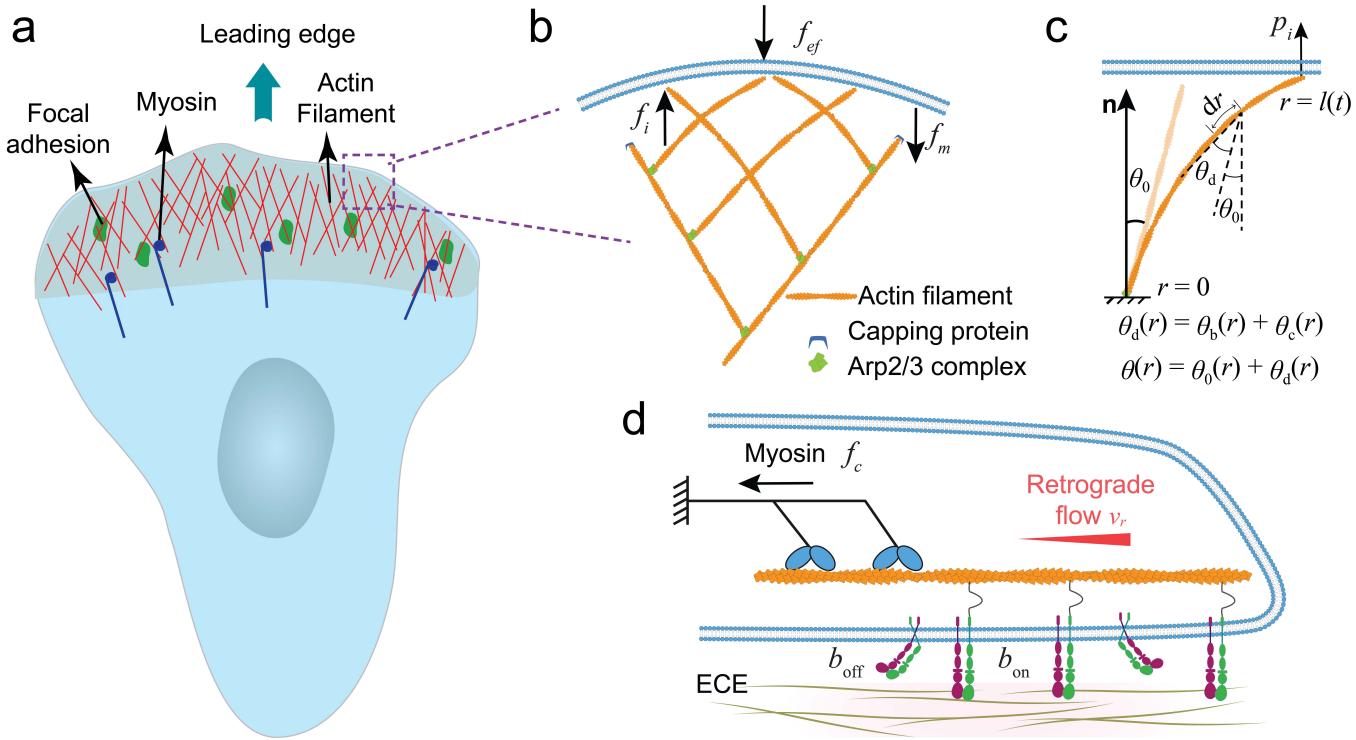


Fig. 1. Schematic representation of AF dynamics during cell migration. (a) Formation of the AF network at the leading edge. (b) Localized view of the branched AF network. The AF deforms as it interacts with the cell membrane, and the deformation influences its subsequent branching behavior. (c) Deformation of an AF and its interaction with the cell membrane. The AF is considered as a beam that can withstand bending and shear deformation. (d) Molecular clutch model with viscous ECE. Myosin contractility pulls on the AFs, leading to retrograde flow, and integrins bound to the ECE provide resistance. AF: actin filament.

$$\sum_i f_i = f_m + f_{ef}, \quad (1)$$

where f_i is the propulsive force of the i th AF. f_m and f_{ef} denote the membrane tension and the resistance of the viscous ECE, respectively. Here, the membrane tension is assumed to be constant during cell migration, and we have also discussed the effect of its variation with ECE viscosity on cell migration (Fig. S5). The resistance from the viscous ECE can be written as a function of viscosity μ and migration speed v

$$f_{ef} = k_1 \mu v, \quad (2)$$

where k_1 is a dimensionless scaling constant.

The polymerization process of AFs is characterized by the addition of individual actin monomers to the barbed ends of an AF. This process is regulated by biological, chemical and mechanical cues^{7,29}. Accordingly, the growth

49 rate of the polymerizing AF is expressed as^{30,31}

$$v_p(t) = \delta [\gamma \cdot g(C_a) \cdot K_p(p_i) \cdot k_{on} - k_{off}], \quad (3)$$

50 where δ is the diameter of an actin monomer, $C_a = C_a(\Phi)$ is the local concentration of actin monomers around
51 the leading edge, which is proportional to the polymerized AF density Φ , $g(C_a)$ denotes the polymerization rate
52 is a function of actin monomer concentration, γ is the diffusion coefficient of actin monomers, k_{on} and k_{off} are the
53 association and dissociation rates of actin monomers to the polymerizing AF, respectively, K_p is the force-dependent
54 probability of AF polymerization. Thus, the length of the polymerizing AF is

$$l(t) = \int_t v_p(t) dt. \quad (4)$$

55 Previous studies indicated that the load-induced curvature of AFs can influence the formation of the branched
56 AF network at the leading edge^{32,33}. Recent cryo-EM experiments demonstrated that the mother filament in contact
57 with the Arp2/3 complex exhibits slight bending and twisting³⁴, as shown in Fig. S1. In this study, the AF is treated
58 as a deformable beam with a circular cross-section. When it interacts with the leading edge membrane, it undergoes
59 bending and shear deformation (Fig. 1c). Thus, the interaction force acting on the i th AF is

$$p_i(t) = \begin{cases} \frac{EI}{\omega(l,t)} \frac{d\theta_b(r,t)}{dr}, & \text{contact} \\ 0, & \text{no contact} \end{cases}, \quad (5)$$

60 where E is the Young's modulus of the AF, $I = \pi d^4/64$ is the second moment of the cross-sectional area of the AF
61 with diameter d , θ_b is the bending angle, and $\omega(l,t)$ is the deflection. The propulsive force is $f_i(t) = p_i(t) \cos \varphi$,
62 where φ denotes the angle between the normal direction of the local leading edge and the direction of cell migration.
63 In addition, the bending force induces conformational changes of the AF, affecting its branching behavior³⁵. The
64 new branch prefers the convex side of a bent filament over the concave side^{30,33}. That is, the Arp2/3 complex, as
65 an actin nucleator, has a higher binding affinity to the convex side of an AF. A bending curvature dependent factor
66 $s_{arp}(\kappa_m)$ is introduced in Eq. 7, which denotes the distance between two adjacent Arp2/3 complex branches along a
67 mother AF, where κ_m is the mean bending curvature of the AF.

68 Moreover, we consider the retrograde flow that occurs in the AFs within the lamellipodia of migrating cells,
69 mainly due to myosin contraction^{36,37}. The pulling force of myosin motors on the AF, driving the actin flow toward
70 the cell center (Figs. 1d and S2). Integrins, which link the actin cytoskeleton to the extracellular environment and
71 provide effective mechanotransduction³⁸, serve as molecular clutches. The retrograde flow speed can be described
72 by the Hill relation^{11,39}:

$$v_r(t) = v_0 \left(1 - \frac{f_{ad}(t)}{f_c} \right), \quad (6)$$

73 where v_0 is the free flow speed, f_c is the myosin stalling force, and f_{ad} is the adhesion force generated by the engaged
74 clutches. Besides, force transmission can lead to strengthening and stabilization of adhesions⁴⁰. In our model,
75 adhesion strengthening is captured by the addition of new clutches when a threshold force acting on the binding
76 clutches is reached (Eq. S10, and Fig. S2). Thus, our model integrates the effects of AF polymerization, branching
77 and retrograde flow, which occur in cell migration. The simulation procedure is described in the Methods section
78 and Fig. S11. The detailed derivation of this mechanobiological model can be found in the Supporting Information.

79 We first calculate the cytoskeleton polymerization rate under different viscous environments, corresponding to
80 different resistances, and compare the obtained resistance-growth rate relation with relevant experiments (Figs. S3
81 and S4). The results show that our model can capture the force-sensitive polymerization behavior, i.e., a slower
82 growth rate under higher resistance conditions, which is consistent with the experimental results²⁶. Our model has
83 accounted for adhesion and actin retrograde flow, which regulate cell motility with cytoskeleton dynamics together.

84 **Increasing ECE viscosity leads to biphasic migration speed**

85 To investigate the effects of increasing ECE viscosity on cancer cell migration, we employ our mechanobiological
86 model to simulate the response of migration speed, focal adhesion dynamics, and retrograde actin flow under varying
87 viscosity conditions. The simulations reveal a biphasic response in the migration speed as ECE viscosity increases
88 (Figs. 2a and S5.). Specifically, the migration speed initially increases, reaching a peak around 700 cP, before
89 decreasing at much higher viscosity. This suggests that moderate increase in viscosity enhances cell motility, which
90 is consistent with experimental observations²². However, beyond this optimal viscosity range, excessive resistance
91 impedes movement, akin to the well-established effects of resistance on cell motility⁴¹. Our model further shows that
92 the adaptive response of a cell is closely linked to changes in its cytoskeletal organization. As viscosity increases, the
93 higher resistance prompts an increase in the number of AFs at the leading edge (Figs. 2b and S3). The increase in
94 filament density and corresponding forces suggests that the cell compensates for the elevated mechanical resistance

95 by recruiting more cytoskeletal elements to maintain sufficient propulsive forces for forward movement. Meanwhile,
96 in a high viscosity environment (e.g., 370 cP), cells exhibit larger focal adhesions (engaged clutches) compared
97 to those in a low viscosity environment (1 cP) (Fig. 2c, left). This indicates an adaptive response wherein cells
98 strengthen their adhesion in high viscosity conditions to sustain propulsive forces during migration. Experimental
99 data on different cells supports this trend, demonstrating an increased number of focal adhesions per cell in high
100 viscosity conditions (Fig. 2c, right;¹⁸).

101 Additionally, our simulations reveal a marked decrease in retrograde flow speed with increasing viscosity,
102 consistent with experimental observations (Fig. 2d). Retrograde flow is significantly slower in high viscosity
103 media (370 cP) compared to low viscosity media (1 cP). This reduction in retrograde flow speed corresponds to a
104 mechanical adaptation of the actin cytoskeleton to heightened adhesions, reflecting a shift in the extent to which the
105 cytoskeleton transmits forces under different conditions. Therefore, the mechanobiological feedback mechanism
106 can be summarized (Fig. 2e). In moderately viscous ECE, the increased resistance enhances cell migration by
107 promoting larger adhesions, denser AF network, and reduced retrograde flow. However, in environments of excessive
108 viscosity, migration speed is ultimately constrained by overwhelming mechanical resistance, which inhibits actin
109 polymerization and diminishes the efficiency of force transmission necessary for cell movement.

110 Deformation of AFs influences the network density

111 Each AF is modeled as a bending beam to explore its dynamic responses to different ECE viscosities. We analyze
112 the changes in AF curvature, density, and the resulting propulsive forces under different viscosity conditions.
113 Simulations reveal that AFs exhibit greater curvatures in high viscosity environments compared to those in low
114 viscosity conditions (Fig. 3a). This increased curvature is a direct consequence of the increased mechanical resistance
115 imposed by the viscous medium, which forces the AFs to bend further. The bending of the AFs facilitates a higher
116 probability of Arp2/3 complex binding to the convex side of the curved AFs, as governed by Eq. (7). As a result,
117 the number of bound Arp2/3 complexes, responsible for nucleation of new AFs at the leading edge, is significantly
118 higher in high viscosity environments (Fig. 3b). This contributes to a denser AF network, which is necessary to
119 generate the propulsive forces required for membrane protrusion under higher resistance conditions. Our results
120 show a strong positive correlation between the curvature of AF and the AF density at the leading edge (Fig. 3c),
121 where increasing AF curvature leads to a significant increase in AF density. This indicates that mechanical bending
122 of AFs directly increases AF density, enabling the cell to adapt to the elevated mechanical resistance.

123 Moreover, the propulsive force generated by the bent AFs is significantly greater in the high viscosity conditions
124 (Fig. 3d). The combination of increased AF curvature and a denser filament network produces a larger force to drive

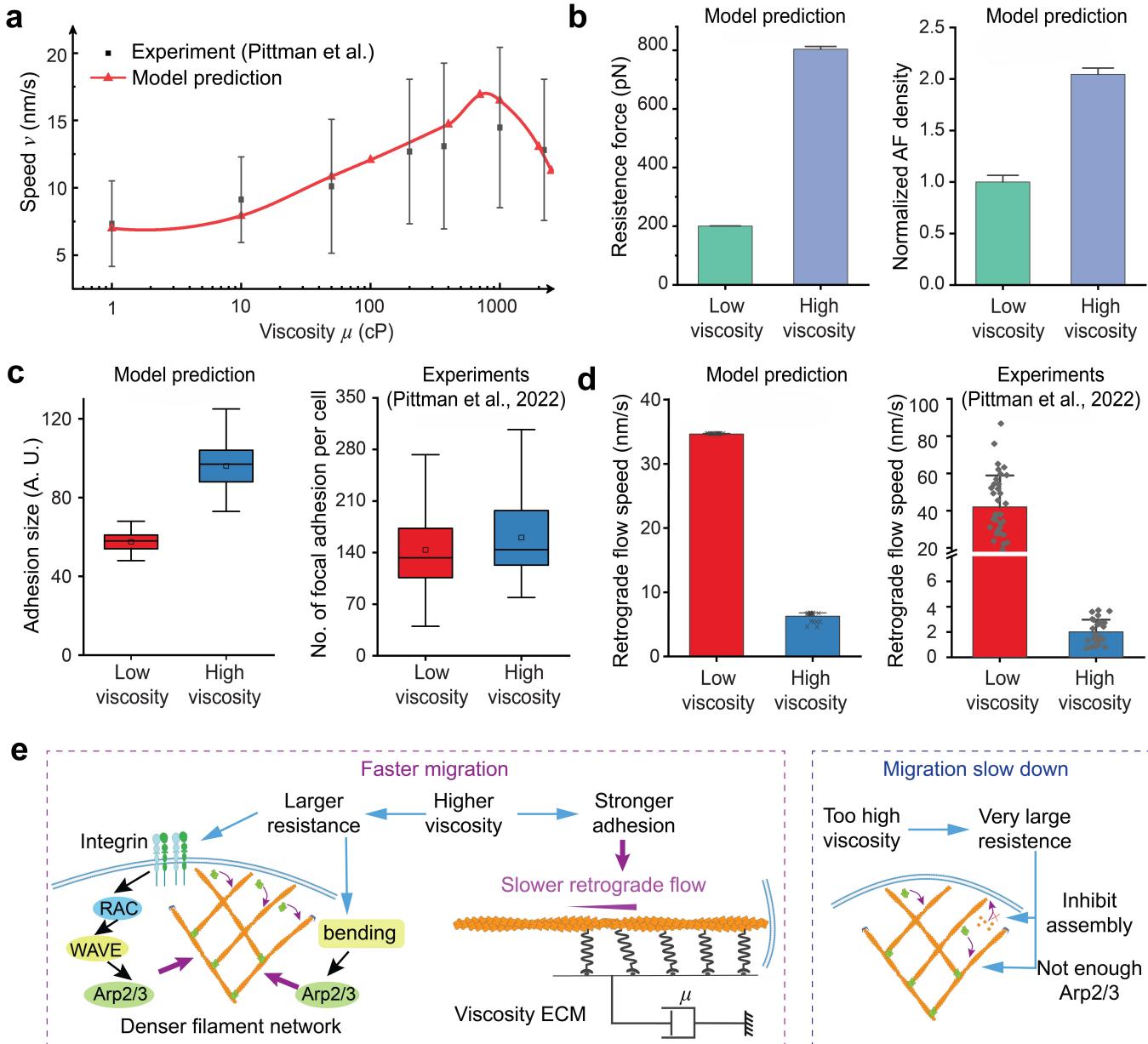


Fig. 2. Differences in the process of cell migration caused by ECE viscosity. (a) Increasing and then decreasing migration speed with increasing ECE viscosity. The simulation results correspond to the average migration speed. The experimental results are the motility of MDA-MB-231 breast cancer cells in viscous media from Pittman et al.²². (b) The model prediction of the resistance force (left) and the AF number (right) at the leading edge. The AF density denotes the number of AFs whose point end is no more than 200 nm from the leading edge. (c) Viscosity induced focal adhesion adaptation. The model prediction indicates that the larger adhesion size in the high viscosity medium (370 cP) compared to the low viscosity medium (1 cP). The adhesion size corresponds to the number of engaged clutches during the simulation. The experiment corresponds to the number of focal adhesions per cell in the low viscosity case (regular medium, 1 cP) and high viscosity case (added 1% high-molecular-weight hydroxypropyl methylcellulose, 370 cP). (d) Effect of viscosity on retrograde flow. The model prediction indicates that the retrograde flow is slower in the high viscosity medium (370 cP) compared to the less viscous medium (1 cP). The experiment corresponds to the same tendency, that is, the retrograde flow is much slower in the high viscosity case (370 cP) than in the low viscosity case (1 cP). (e) Mechanobiological mechanism of ECE viscosity on cell motility. The migration speed corresponds to the statistics after the simulation is stabilized.

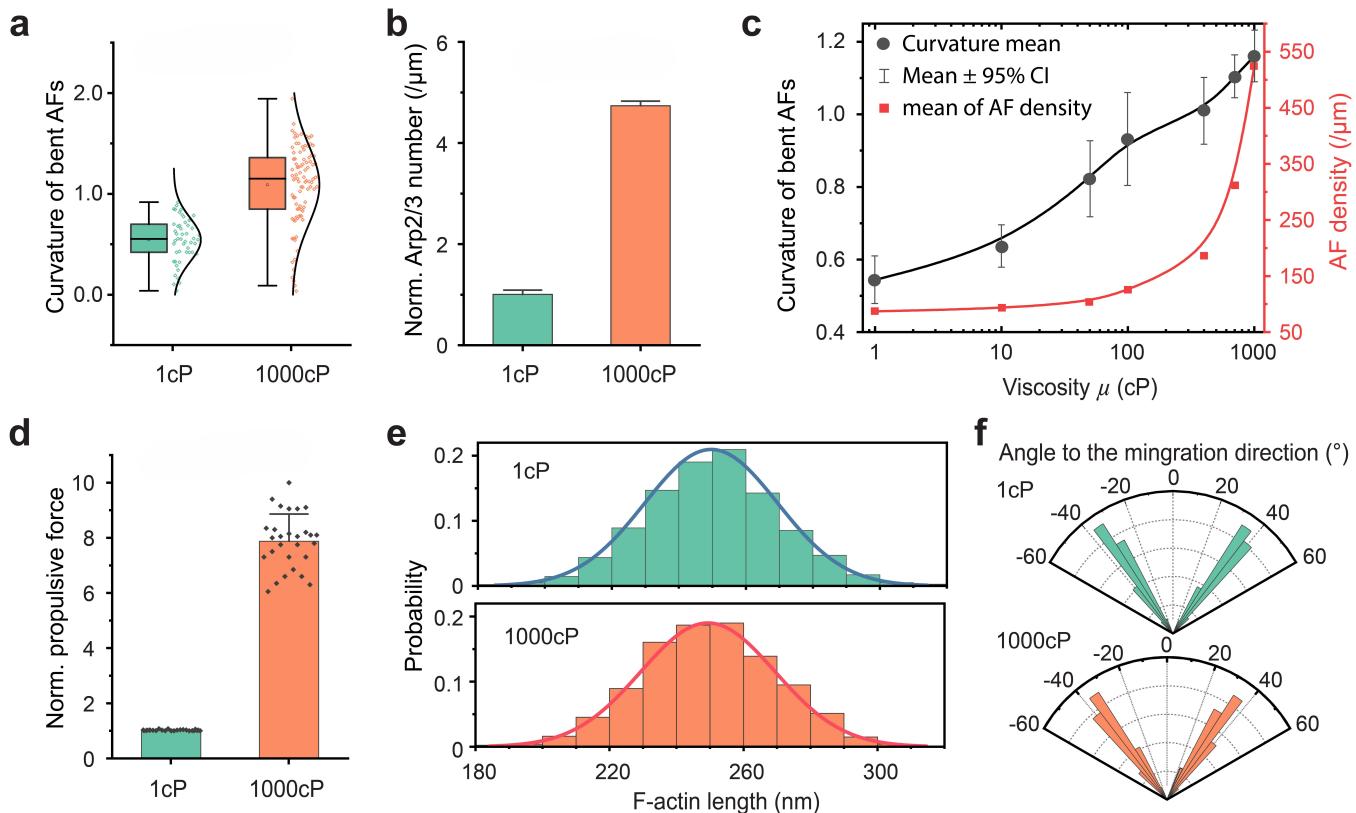


Fig. 3. Deformation and distribution of AFs. (a) The mean curvature of bent AFs in low and high viscosity ECE. The curvature is much larger in the latter condition. (b) The normalized number of Arp2/3 bound to the mother AFs. More bound Arp2/3 is related to more newly generated daughter AFs, that is, a denser AF network. Error bars, s.d. (c) The relation between the curvature of the bent AF and its density at the leading edge. The AF density is proportional to the statistical mean curvature. (d) The propulsive force attributed by bent AFs. The denser AF network and larger curvature in the high viscosity ECE lead to a larger propulsive force. Error bars, s.d. (e) The length of an individual AF follows a Gaussian distribution, and (f) at an approximate angle of $\pm 35^\circ$ from the forward direction of migration.

125 membrane protrusion, a critical adaptation to overcome the mechanical challenges of high viscosity environments.
126 The lengths of individual AFs follow a Gaussian distribution in all conditions (Fig. 3e), and the initial direction
127 before bending is aligned within an angular range of approximately $\pm 35^\circ$ from the forward migration direction
128 (Fig. 3f). Additionally, filament length also influences cellular and molecular behaviors, with longer filaments
129 supporting increased cell migration and contributing to a denser AF network (Fig. S6). These results highlight the
130 critical interplay between filament geometry and network density in regulating the cellular response to mechanical
131 resistance. Such insights offer a deeper understanding of how cancer cells adapt their cytoskeletal structure to
132 optimize migration in challenging microenvironments.

133 **Actin monomer concentration influences migration speed**

134 The availability of profilin-bound actin monomers significantly affects AF polymerization. Our simulations show
135 that as ECE viscosity increases, the actin monomer concentration near the leading edge is upregulated to sustain
136 AF polymerization (Fig. 4a). Within a given viscosity condition, the concentration of actin monomers affects AF
137 polymerization and then acts as a change in migration speed (Fig. 4b). Increased availability of actin monomers
138 enhances the polymerization of AFs, which provides propulsive forces necessary for efficient movement, thus
139 leading to faster migration speeds. To further elucidate the role of actin monomer availability, we examined three
140 specific cases (Fig. 4c). Normally, the monomer concentration is below $22.5\mu\text{M}$ under low viscosity conditions
141 (1 cP and 100 cP), as shown in Fig. 4a. Increasing the monomer concentration to $22.5\mu\text{M}$ represents an elevated
142 supply. While in the high viscosity condition (1000 cP), the monomer concentration normally exceeds $22.5\mu\text{M}$.
143 Reducing the available monomer concentration to $22.5\mu\text{M}$ in this setting limits the availability of actin monomer.
144 Fig. 4d, illustrating the migration speeds, shows that an increased supply of actin monomers significantly boosts the
145 migration speed, whereas a reduced availability of actin monomers significantly slows the migration.

146 The corresponding AF density at the leading edge (Fig. 4e) shows that the AF density is higher in the low
147 viscosity conditions (1 cP and 100 cP) with elevated actin monomer supply ($22.5\mu\text{M}$) compared to the normal
148 condition (Fig. 3c). In contrast, reducing the monomer supply to $22.5\mu\text{M}$ in the high viscosity condition (1000 cP)
149 results in a lower AF density compared to the normal condition. These AF adjustments correspond to changes in
150 cell motility. Then, we present a phase diagram that maps the interplay between ECE viscosity and actin monomer
151 availability, delineating the optimal conditions for maximizing cell migration speed (Fig. 4f). It shows that low
152 monomer concentrations or excessively high viscosity can impair movement, while an optimal combination of
153 these factors supports rapid migration. These results emphasize that effective cell migration depends on a tuned
154 balance between the availability of associated proteins and the mechanical resistance provided by the extracellular

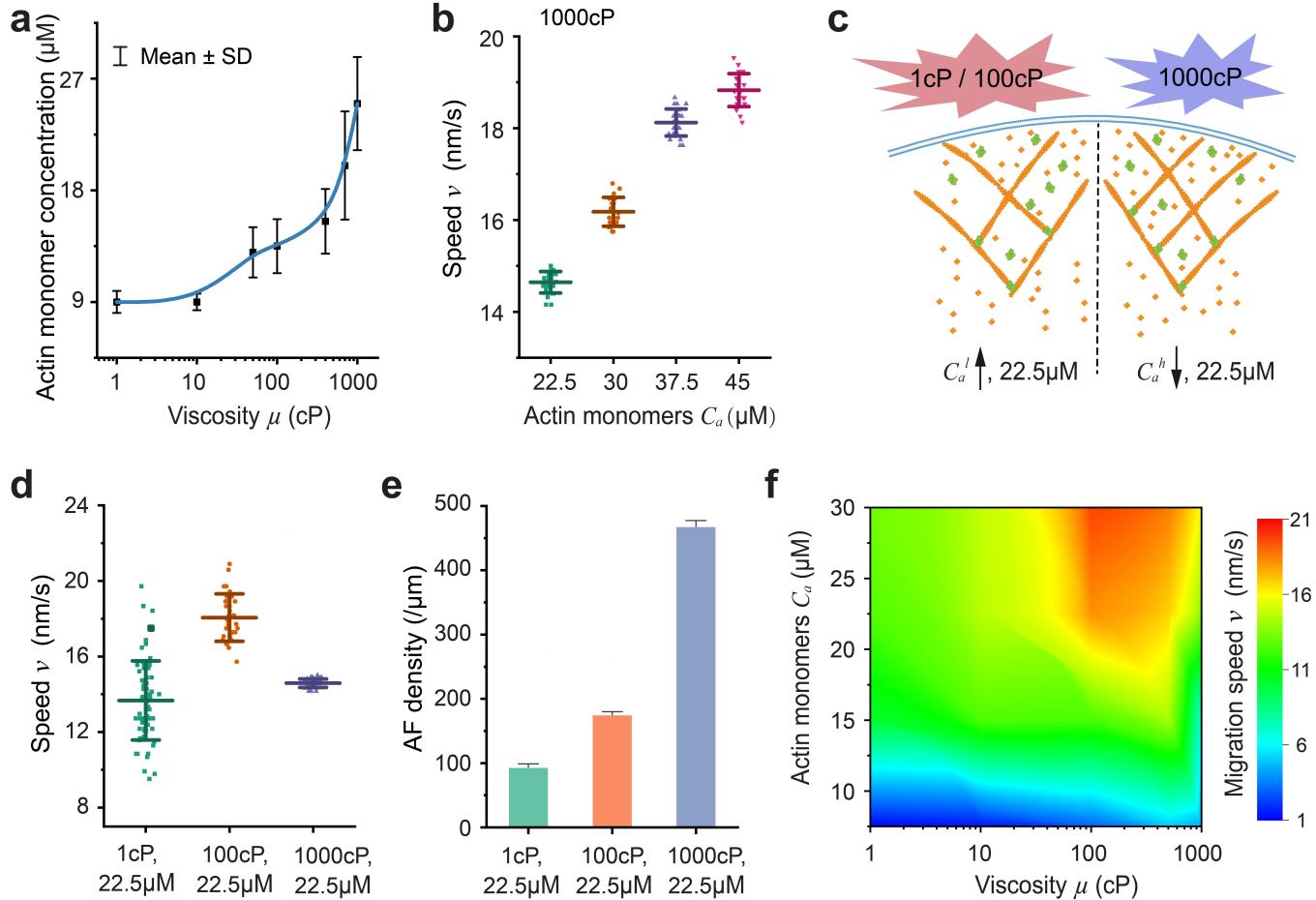


Fig. 4. Influence of actin monomers and viscosity on the AF network and cell migration. (a) Actin monomer concentration at each step for different viscosity environments. Error bars, s.d. (b) Influence of actin monomer concentration on migration speed. Higher concentration promotes AF polymerization, leading to higher migration speed. (c) Cases simulating the influence of actin monomers on cell behavior. For lower viscosity cases (1 cP/100 cP), the actin monomer concentration is less than $22.5\mu\text{M}$ in normal situation, thus, the concentration $22.5\mu\text{M}$ means increasing the supply of actin monomers. In contrast, the concentration $22.5\mu\text{M}$ means reducing the supply of actin monomers for the high viscosity case (1000 cP). (d) Migration speed of the three cases presented in (c). (e) The AF density at the leading edge of the three cases. (f) Phase diagram of biochemical cues for cell migration speed, which shows the optimal ECE viscosity and actin monomer concentration for cell migration.

155 environment, offering insights into how cancer cells adapt their motility to varying physical conditions.

156 **Integrin ligand density and myosin contraction influence retrograde flow**

157 Integrin ligand density and myosin concentration play critical roles in modulating retrograde flow and cell motility
158 through mechanobiological interactions such as cell-ECE adhesion, integrin activation, and myosin-induced con-
159 tractility^{42–44}. Our results show that cell-ECE adhesion becomes stronger with increasing ECE viscosity (Fig. 5a,
160 left). This increased adhesion contributes to a reduction in retrograde flow speed under higher viscosity conditions
161 (Fig. 5a, right). This slowed retrograde flow enables the generation of more effective propulsive forces to drive
162 forward movement. We further investigate how ligand density, which refers to the availability of binding sites such
163 as integrins or adhesion complexes⁴⁵, influences retrograde flow speed and migration speed. When ligand density
164 is low, retrograde flow speed is higher, indicating weaker cell-ECE adhesion and allowing the actin network to
165 move backwards more rapidly (Fig. 5b, left). This condition corresponds to the reduced migration speed (5b, right).
166 However, as the ligand density increases, the retrograde flow decreases, and the cell migration speed increases due
167 to the enhanced propulsion of AF networks. Interestingly, beyond a certain threshold, further increases in ligand
168 density have little effect on migration speed (5b, right), indicating a plateau where adhesion strength is maximized
169 and additional binding sites do not significantly enhance motility. This phenomenon is in partial agreement with
170 experimental results⁴⁶, while the biphasic dependence of cell migration speed on ligand density^{46,47} requires further
171 modeling investigation.

172 Myosin contractile forces can also affect retrograde flow and cell migration speed. In high viscosity ECM,
173 increased myosin contraction leads to slightly faster retrograde flow (Fig. 5c), and this change in retrograde flow has
174 a small effect on migration speed (Fig. 5c). That is, myosin contractility has little influence on migration speed in
175 the high viscosity case, which is consistent with some experimental observations^{22,48}. Besides, increased myosin
176 contractility within a range can strengthen adhesion (Fig. 5d), because of the enhanced integrin binding mechanism⁴⁰.
177 As myosin contractility increases further, adhesion may become weaker (Fig. S7a). We summarize these relations
178 with a schematic representation (Fig. 5e), illustrating how ligand density and myosin contraction collectively
179 influence retrograde flow and migration speed. Higher ligand density enhances adhesions, slows retrograde flow, and
180 promotes forward migration. On the other hand, elevated myosin contraction pulls on the AF network, accelerating
181 retrograde flow and, as a result, slowing forward migration. While in the high viscosity ECM, the effect of myosin
182 contractility has less effect because retrograde flow is very slow as adhesion strengthens, and AF polymerization is
183 the main protrusion factor (Fig. S7b).

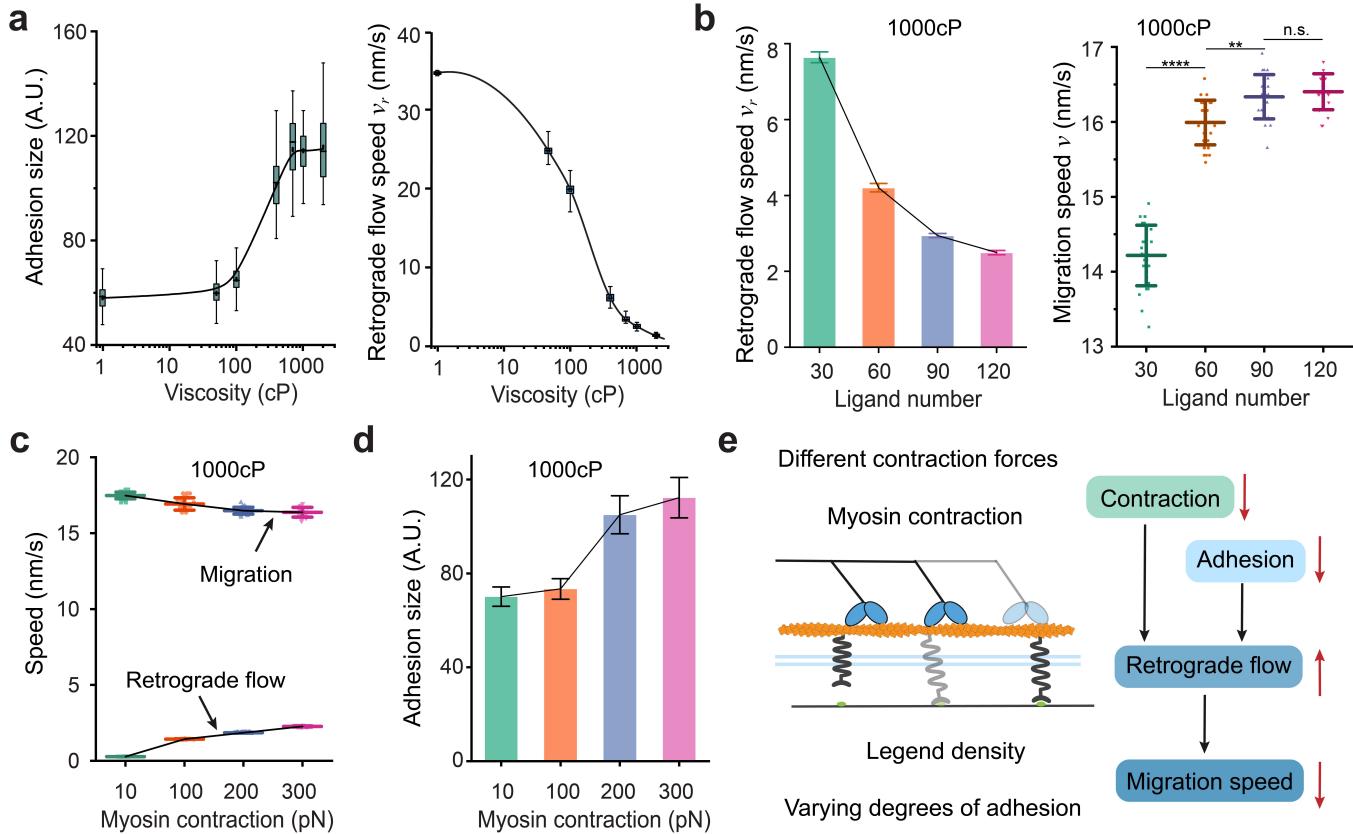


Fig. 5. Mechanobiological cues for actin flow and cell migration. (a) Relation between ECE viscosity and cell-ECE adhesion (left) or retrograde flow speed (right). Whisker range 1 ~ 99. (b) Relation between retrograde flow speed and the number of ligands for the high viscosity case (1000cP). Fewer ligands lead to fast actin flow (left), and correspondingly slower migration speed. When the ligand becomes dense enough, it has less influence on the migration speed. Error bars, s.d. (c) Effect of myosin contraction on retrograde flow and migration speed. Greater myosin contraction leads to slightly higher retrograde flow speed and slower migration speed. Error bars, s.d. (d) Effect of myosin contraction on adhesion. Within a certain range, increased myosin contraction leads to strengthened adhesion. (e) Schematic mechanism of ligand density and myosin contraction for actin flow and migration.

184 **Change in viscosity leads to mechanobiological migration memory**

185 Cell migration memory is an emergent property that allows cells to utilize past experience to inform future movements,
186 with significant impacts for pathological processes^{24,25}. The cytoskeleton-based mechanical memory was reported in
187 some experiments^{49,50}, and might be related to the formation of long-term memory⁵¹. In this study, we elucidate the
188 mechano-chemical coupled mechanisms of short-term cell migration memory. When a cell is transferred from the
189 low viscosity ECE (1 cP) to high viscosity ECE (1000 cP), its migration speed initially drops sharply, indicating the
190 immediate mechanical impact of the higher viscosity (Fig. 6a). As the cytoskeleton adapts by polymerizing and
191 forming a denser AF network (Fig. 6c), the migration speed gradually increases. By this mechanism, the cell can
192 regain its speed despite increased environmental resistance. Fig. 6b shows more distinctly that the cell moving in
193 the low viscosity environment exhibits slow migration speeds, but it achieves faster movement through cytoskeletal
194 adaptation when transferred to the high viscosity condition. This increase in migration speed under high viscosity
195 conditions suggests that the cell requires sufficient propulsive force to overcome the added resistance. Our results
196 show that the density of the AF network increases with rising viscosity (Fig. 6c), indicating that cells adapt to higher
197 mechanical resistance by constructing a denser cytoskeletal structure, enabling them to exert stronger forces against
198 the high viscosity ECE.

199 Interestingly, in the reverse case, when a cell is transferred from the high viscosity (1000 cP) to low viscosity (1
200 cP) condition, the migration speed first increases rapidly and then decreases. However, compared with the speed for
201 cells consistently exposed to low viscosity, the migration rate stabilizes at a higher level, similar to that observed
202 in the high viscosity condition (Fig. 6d). Fig. 6e compares the migration speeds in both constant and changing
203 viscosity scenarios. The cell transitioning from high to low viscosity can maintain a higher migration speed than
204 that under constant low viscosity conditions. This phenomenon suggests a short-term migration memory, wherein
205 the cell can "remember" the motility developed in the high viscosity environment, allowing it to maintain faster
206 movement. By interacting with epigenetic changes, this cytoskeleton-based memory may favor long-term memory
207 formation²³.

208 To further elucidate this memory effect, we examine the AF density when a cell is transferred from high to
209 low viscosity. Although the AF density decreases upon transition to the low viscosity environment (Fig. 6f), it
210 still remains higher than that of the cell always exposed to low viscosity conditions (Fig. 6g). Additionally, the
211 cell transferred from high to low viscosity encounters greater resistance than that in continuously low viscosity
212 environments due to their higher migration speed (Fig. 6h), which favors AF branching. This residual cytoskeletal
213 density, coupled with elevated actin monomer concentration near the leading edge (Fig. S8), induces the increased

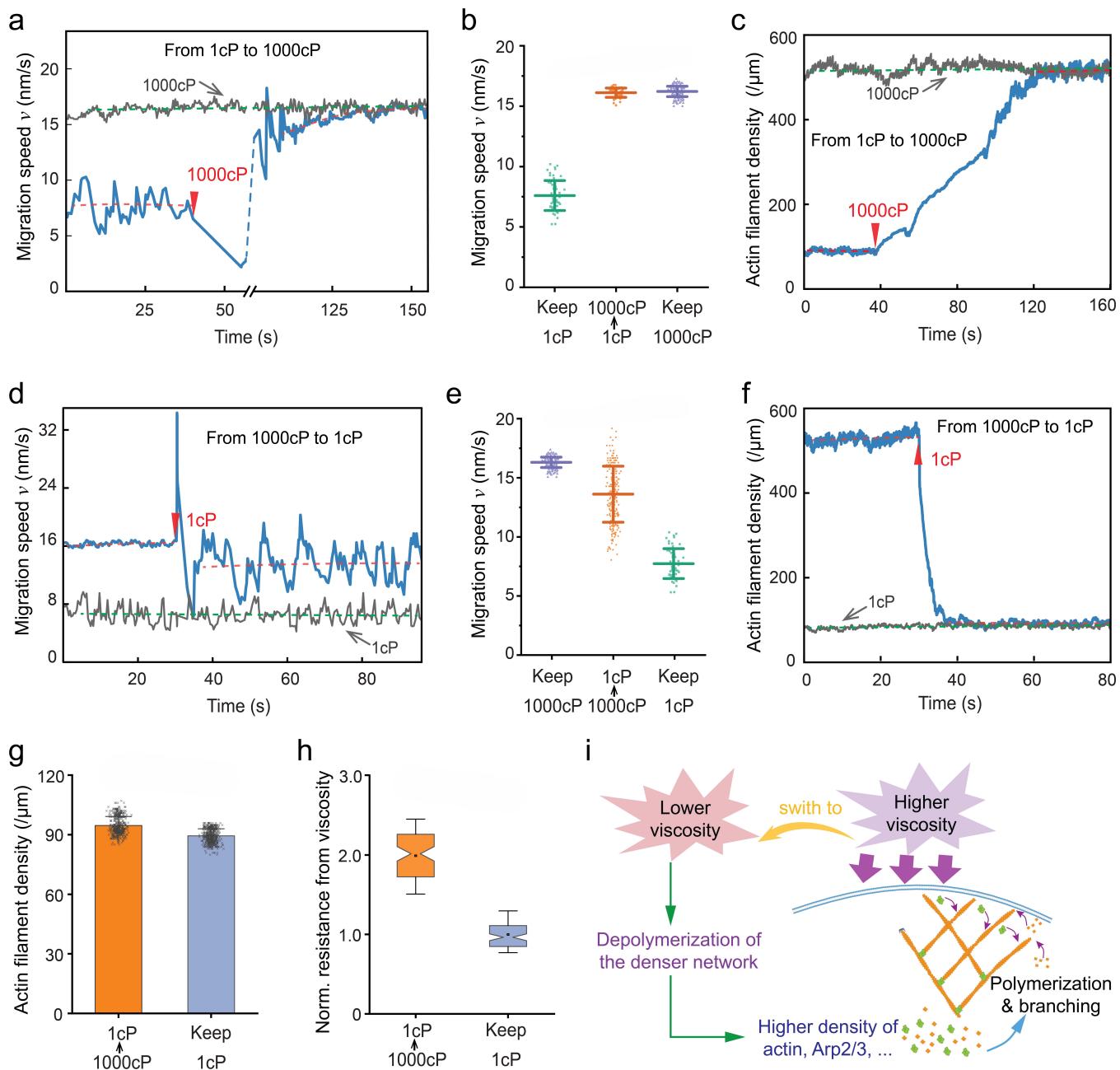


Fig. 6. Changes in the viscosity of ECE and the effect on cell migration and cytoskeleton evolution. (a) The change in migration speed when a cell is transferred from low to high viscosity. (b) Comparison of migration speed in the viscosity changing and invariant cases. Error bars, s.d. (c) The change of AF density with viscosity. The AF density becomes larger with increasing viscosity. (d) The change in migration speed when a cell is transferred from high to low viscosity, with a sufficient actin monomer supply. (e) Comparison of migration speed in the viscosity changing and invariant cases. The migration speed retains some memory, that is, the cell maintains a similar speed in the high viscosity case after being transferred to the low viscosity case. Error bars, s.d. (f) The AF density changes as ECE becomes more viscous. (g) The difference in AF density for the case of transfer from high viscosity and invariant low viscosity. The small larger density is significant for the cell to maintain higher speed, which can provide greater propulsive force (h). (i) The short-term memory mechanism of cell migration. Cells in a high viscosity environment are stimulated to synthesize more relevant proteins, such as Arp2/3 and actin, to form a denser AF network. When a cell is transferred to a low viscosity environment, the cell still maintains the depolymerization of the dense AF network far from the leading edge, to produce a higher protein density, which is conducive to cytoskeletal polymerization. This in turn generates a greater propulsive force for cell movement. The time zero is calculated from the time after the AF network has produced stable propulsion. All traces in a, c, d and f are the trace after steady migration.

214 migration speed in a low viscosity ECE. AF polymerization is influenced by actin monomer concentration near the
215 leading edge, which is supplied from the cytosolic pool, and by recycling of depolymerized filaments (Fig. S8).
216 Additionally, the enhanced contractility at the cell rear under high viscosity²³ may further drive forward flow of
217 actin monomer in the cytosolic pool to the leading edge⁵². In addition, the change and different fluctuations in
218 actin monomer concentration can affect the duration of migration memory (Methods and Fig. S9). To explain this
219 memory phenomenon, we propose a mechano-chemo-biological mechanism (Fig. 6i). When cells are exposed
220 to a high viscosity environment, they increase the assembly of AF network that facilitates movement in resistant
221 conditions. Even after transitioning to a lower-viscosity environment, this enhanced network assembly allows the
222 cells to maintain a higher migration speed than those that were never exposed to high viscosity. This mechanical
223 memory may be further supported by potential epigenetic changes within the cell (Fig. S10b), which requires further
224 in-depth investigation.

225 Discussion

226 Our model demonstrates a linkage between AF deformation and AF network formation through mechano-chemo-
227 biological mechanisms. Together with the consideration of retrograde flow, the model reveals a biphasic response
228 in migration speed as ECE viscosity increases (Figs. 2a and S5). This model shows that increased viscosity can
229 strengthen adhesions, reduce retrograde flow, and modulate actin polymerization dynamics, allowing the cells to
230 sustain faster migration speeds up to a threshold where resistance becomes too high. This mechanistic insight
231 underscores how adhesion strength and actin dynamics collectively mediate viscosity-dependent cell motility. The
232 phenomenon, wherein cancer cells move faster in more viscous ECE, is consistent with experimental observa-
233 tions^{20,22,23}. The biphasic pattern indicates that cytoskeletal adaptability is not only a function of the related proteins
234 availability but is also tuned by the physical properties of the extracellular environment.

235 Actin polymerization is an essential process to generate propulsive forces for cell movement. The actin
236 monomer availability is a key factor in driving AF polymerization. In low viscosity environments, cells require less
237 actin polymerization to overcome low mechanical resistance, while in high viscosity environments, the demand
238 for actin monomers increases as cells must exert greater force to migrate through a denser and more resistant
239 medium. The efficient actin monomer recycling and transport ability^{52,54–56} helps maintain the local actin monomer
240 concentration near the leading edge at a sufficient high level to guarantee continuous polymerization. Besides,
241 elevated actin monomer concentrations promote rapid polymerization, translating to faster migration speeds (Fig.
242 4b). Biomechanical factors such as ligand density and myosin contractility also affect cell migration. Ligand
243 density, which correlates with cell-ECE adhesion dynamics via integrins, can influence the adhesion strength,

Table 1. Key cellular and subcellular behaviors and predictive capabilities of our model

	Experimental phenomena	Our model
AF Polymerization	Slower polymerization for larger resistance ²⁶	Yes (Fig. S4)
	Denser AF network for larger resistance ^{26,29}	Yes (Figs. 3c, S3)
Adhesion	Strengthened adhesion with increasing force applied to integrins ^{40,53}	Yes (Figs. 5d, S2)
	Increasing viscosity strengthens adhesion ^{18,23}	Yes (Fig. 5a)
	Increasing viscosity slows retrograde flow ^{22,23}	Yes (Fig. 5a)
Cell motility	Actin but not myosin dominates cell motility in the high viscosity condition ^{22,48}	Yes (Figs. 5c, S7b)
	Increasing viscosity speeds up migration ^{18,20,22,23}	Yes (Figs. 2a, S5)
	Migration slows as viscosity becomes too high (not enough experiments)	Yes (Figs. 2a, S5)
Memory effect	Cell migration speed remains high in low viscosity environment after pretreatment in high viscosity environment (long-term) ²³	Long-term is not included
	AF polymerization rate becomes faster in low resistance after pretreatment in high resistance condition (short-term) ⁴⁹	Yes (Figs. 6d, 6e, S9)

244 thereby impacting retrograde flow. Optimal ligand density facilitates effective adhesion, promoting motility. Myosin
245 contractility also plays a central role in regulating retrograde flow. Increased myosin activity enhances cytoskeletal
246 tension, driving retrograde flow. However, in high viscosity environments, retrograde flow is naturally slow, and
247 the impact of myosin contractility can be diminished. This is verified by the experimental results that inhibition of
248 myosin has a negligible effect on cell motility in high viscosity ECE²². This nuanced interaction between adhesion,
249 contractility, and viscosity underscores the complicated mechanism, where cells regulate cytoskeletal tension and
250 adhesion to achieve optimal migration.

251 Finally, this study reveals the mechanisms of the cytoskeleton-based short-term migration memory, wherein
252 cells exposed to a high viscosity environment can maintain faster migration speeds even after transitioning to low
253 viscosity conditions. This phenomenon suggests that cells can store information from their previous mechanical
254 environment and use it to maintain enhanced motility. At the molecular level, this memory is likely related to the
255 increased concentration of cytoskeletal proteins, such as actin monomers and Arp2/3 complexes, which is primarily
256 due to rapid monomer transport and efficient actin recycling mechanisms (Figs. 6i and S8). The persistence of
257 these proteins results in a denser AF network, providing sustained structural integrity and enabling the cell to
258 move more quickly. A biological implication of this memory effect is particularly relevant in the context of cancer
259 metastasis. During metastasis, cancer cells need navigate through diverse tissue environments, ranging from high
260 viscous extracellular matrices to more fluid environments like blood vessels or lymphatics. Maintaining enhanced

261 motility after transitioning from dense to softer tissues might be a significant ability of cancer cells, allowing them
262 to more effectively disseminate and colonize distant sites. Our model demonstrates how biomechanical cues can
263 drive short-term cytoskeletal memory after viscosity transitions, which is similar to the experimentally observed
264 mechanical memory^{49,50}. The biochemical pathway-mediated transcriptional changes^{23,57} may explain the long-
265 term regulatory mechanisms for migration behavior. To fully integrate those mechanotransduction pathways and
266 transcriptional feedback mechanisms^{24,58}, future studies are needed to develop whole-cell models that encompass
267 both cytoskeletal dynamics and transcriptional changes, including their interaction⁵¹. Such models would provide a
268 unified understanding of how mechanical and biochemical cues synergistically regulate migration behavior across
269 different time scales.

270 In summary, we have highlighted the effects of biomechanical cues, particularly ECE viscosity, on modulating
271 cytoskeletal behavior, adhesion and cell motility, as shown in Table 1. We reveal how a cell dynamically integrates
272 mechanical resistance with biochemical signals to tune its cytoskeletal structure and migration speed. Furthermore,
273 the discovery of cytoskeleton-based migration memory introduces a new layer of complexity to our understanding
274 of cellular adaptation, revealing how transient changes in the mechanical environment can induce effects on the
275 migratory machinery. This work provides potential therapeutic insights for inhibiting abnormal cancer cell migration,
276 suggesting that manipulating tissue mechanics or targeting the molecular pathways involved in AF polymerization
277 and migration memory may offer novel strategies for limiting metastatic progression.

278 Methods

279 Polymerization of the actin filament

280 AF polymerization is a dynamic process corresponding to the addition of actin monomers primarily at the barbed end
281 and the removal of monomers at the pointed end. In the simulation, we assume that the process involves the addition
282 of actin monomers only at the barbed end while the removal of actin monomers occurs at the pointed end. AFs
283 polymerize over time and may be capped by capping proteins or interact with the membrane. The polymerization
284 rate can be obtained from Eq. (3). The concentration of actin monomers C_a has an initial minimum value C_a^0 . This
285 concentration in the following simulation is related to the density of the AF network already far from the leading
286 edge because the AF depolymerization can release actin monomers, which supplies free actin monomers^{55,56,59,60}.
287 Additionally, there is an upper limit to the concentration C_a^c due to the diffusion limit of actin monomers. That
288 is, the temporal concentration of actin monomers is determined by the AF density and the diffusion coefficient,
289 corresponding to Φ and γ in Eq. (3), respectively. The force-dependent probability of AF polymerization is an
290 exponential relation as in Eq. (S1). The association rate k_{on} and dissociation rate k_{off} of actin monomers to the

291 polymerizing AF are constant during the simulation.

292 As the filaments polymerize, they might interact with the membrane. Branching may occur on the convex side
293 of the mother filaments and new daughter filaments begin to polymerize at the Arp2/3 binding site^{30,33}. The number
294 of daughter filaments of a mother filament can be determined according to

$$s_{\text{arp}}(\kappa_m) = \begin{cases} 252, & 0 \leq \kappa_m < 0.3 \\ 144, & 0.3 \leq \kappa_m < 0.6, \\ 90, & \kappa_m \geq 0.6 \end{cases} \quad (7)$$

295 where s_{arp} is the distance between two Arp2/3 binding to a mother filament. Thus, the number of Arp2/3 binding to
296 the filament is $N_{\text{arp}} = \lfloor l/s_{\text{arp}} \rfloor$, where $\lfloor \cdot \rfloor$ denotes the rounding down operator. Additionally, the critical curvature is
297 set because the experimental observation shows that the bending angle is several degrees³⁴. Besides, different types
298 of cells might have different abilities to generate branched network, different s_{arp} can simulate these differences.
299 Experimental results have shown that the angle between the branched daughter filament and mother filament is about
300 70° ^{7,61}. During the simulation, the branch angle is set as a random value between $[68^\circ, 72^\circ]$.

301 **Retrograde flow**

302 The polymerization of AF is also accompanied with the retrograde flow, which is simulated by the molecular clutch
303 model in this study. The retrograde flow speed of the whole network is simply calculated using Eq. (6). The
304 adhesion, corresponding to clutches in the molecular clutch model, is regulated by the ECE viscosity and the force
305 acted on the clutch (Eq. (S14)). The adhesion strengthening is implemented by adding new clutches with a dynamic
306 rate (Eq. (S10)). Then, the engaged clutch number is updated using Eq. (S8 - S10). Hence, the retrograde flow speed
307 at the next time step can be updated using Eq. (6). The differences between the molecular clutch model used in this
308 study and other developed related models are listed in Table S1, and the detailed formulation of the molecular clutch
309 model can be found in the Supporting Information.

310 **Simulation process**

311 At the start of the simulation, the starting points of the mother filaments are initialized. First, we select a 2-
312 dimensional region with a width of $x \in (-500, 500\text{nm})$, assuming that the cell membrane is flat within this width
313 range. The initial position of the cell membrane is $y = 60\text{nm}$. The polymerization starting points of the initial
314 mother filaments are randomly generated and distributed in the height range $y \in (0, 20\text{nm})$ and the width range
315 $x \in (-500, 500\text{nm})$. The number of initial mother filaments is N_{mf} . To ensure that the number of mother filaments

316 is roughly uniformly distributed over the width, the generation was done by dividing the width into 10 parts, and
317 randomly generating $N_{\text{mf}}/10$ coordinates of starting points in each of these parts. The orientation of AFs relative to
318 the cell migration direction is generally around $\pm 35^\circ$ ^{62,63}. Thus, the initial angle of the generated mother filaments
319 satisfies a Gaussian distribution $\phi \sim N(\pm 35^\circ, 5^\circ)$. Next, the AF polymerizes by adding individual actin monomers
320 whose radius is $\delta = 2.5\text{nm}$ before growing to its maximum length or being capped, and the filament closer to the
321 leading edge has a higher priority for assembly with actin monomer. The capping situation includes two scenarios,
322 one is the randomness of generated length (a Gaussian distribution $N(l, \pm 20\text{nm})$) and the other is the stopping
323 of polymerization beyond a certain distance from the membrane. The polymerization rate is calculated using Eq.
324 (3), where polymerization slows as actin monomer decreases. In the simulation, the initial polymerization rate
325 is expressed as $v_p = \delta [\gamma C_a^0 K_p(p_i) k_{\text{on}} - k_{\text{off}}]$. At each time step, the number of actin monomers is calculated by
326 $N_a(t) = \min[\max(C_a^0 * 4\mu\text{M}^{-1}, A_{\text{re}} + A_{\text{cp}}), C_a^c * 4\mu\text{M}^{-1}]$, where A_{re} equals to the number of aged AFs divided by a
327 random number $\mathcal{R}(4, 6)$, i.e. the actin monomers from recycling. The pointed ends of the aged AFs are between
328 1000 nm and 1200 nm from the leading edge, i.e., actin filament depolymerization is assumed to occur one micron
329 from the leading edge. For simplicity, the actin monomer from the cytosolic pool is set as $A_{\text{cp}} = A_{\text{re}}$. When these
330 actin monomers are depleted, the simulation proceeds to the next time step. Meanwhile, there is a retrograde flow
331 whose speed is calculated by Eq. (6). The retrograde distance is $\Delta s_r = v_r \Delta t$ at each time step, where $\Delta t = 0.02\text{ s}$ is
332 the incremental time. When the polymerizing AF contacts the membrane, it bends. There is an interaction force
333 calculated by Eq. (5), which acts as the driving force for cell migration. When the propulsive forces of all AFs
334 (left side of Eq. (1)) are greater than the resistance forces (right side of Eq. (1)), the cell migrates with a step size
335 of Δd . The migration speed is $v = \Delta d / (t_j - t_{j-1})$, where $(t_j - t_{j-1})$ is the time from last to current migration step.
336 Meanwhile, after bending of the AF, the branching phenomenon, that is, the binding of Arp2/3, will occur according
337 to Eq. (7). Besides, the interaction force between the AF and the membrane will inhibit the polymerization speed in
338 the next time step, as described by the function $K_p(p_i)$. In summary, the flowchart of this simulation process can be
339 found in Supporting Information Fig. S11.

340 To simulate the memory effect, the actin monomer concentration remains high ($C_a \simeq C_a^c$ in Fig. 6 and different
341 fluctuations in Fig. S9) near the leading edge after the cell has transferred from the high to the low viscosity
342 condition, based on the mechanism shown in Fig. S8. The larger resistance in the high viscosity ECE leads
343 to denser F-actin networks to push the membrane forward^{26,29}. The branched Arp2/3–actin filament network
344 subsequently leaves the ‘activation zone’ as the membrane is pushed forward. Then the aged AFs undergo
345 debranching and depolymerization, and the actin monomers released by the AF network disassembly can be reused
346 for subsequent rounds of polymerization^{55,56,59,60}. The actin monomer transportation can also support continuous

347 polymerization^{52,54}.

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Author contributions statement

Z.L., X.C. and X.-Q.F. designed research; Z.L., X.C. and X.-Q.F. developed computational framework; and Z.L., X.C. and X.-Q.F. analyzed data and wrote the paper.

Competing interests

The authors declare no competing interests.